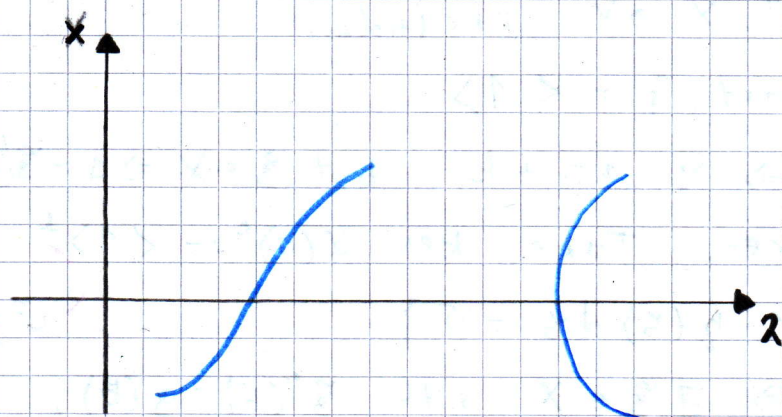


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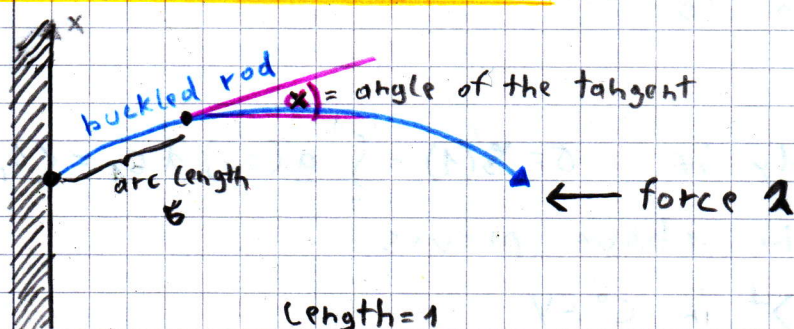
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$\dot{x} = f(\lambda, x)$, $x \in \mathbb{R}^N$, center manifold $\Rightarrow$ exchange of stability

Example 1: Euler's rod

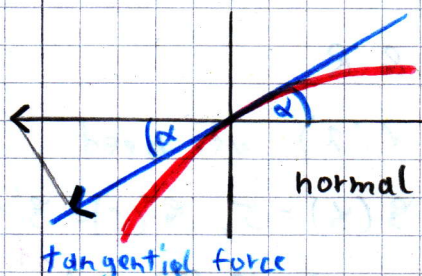


elasticity model:

$$x''(s) \sim -\lambda \sin x(s)$$

$$x''(s) + \lambda \sin x(s) = 0$$

$$x'(0) = x'(1) = 0$$



normal force $= -\lambda \sin \alpha$

tangential force

How to rewrite $\otimes$ as $f(\lambda, x) = 0$

$$x(\cdot) \in X := C^2([0, 1], \mathbb{R}) \cap \{x'(0) = x'(1) = 0\}$$

$$Y := C^0([0, 1], \mathbb{R})$$

Then $f: \mathbb{R} \times X \rightarrow Y$, C^∞

$$\text{for } f(\lambda, x) := x''(\cdot) + \lambda \sin x(\cdot)$$

$$\text{Note: } f(\lambda, 0) = 0$$

$$\mathcal{L}(\lambda) \mathfrak{z} = \mathbb{D}_x f(\lambda, 0) \mathfrak{z} = \mathfrak{z}''(\cdot) + \lambda \mathfrak{z}(\cdot)$$

$\mathcal{L}(\lambda): X \rightarrow Y$ is linear

$$\mathcal{L}(\lambda) = \left(\frac{\partial}{\partial s}\right)^2 + \lambda \text{id}$$

compact!

Arzelà-Ascoli
from C^2 to C^0

Therefore $\mathcal{L}(\lambda)$ Fredholm $\Leftrightarrow \mathcal{L}(0) = \left(\frac{\partial}{\partial \epsilon}\right)^2$ Fredholm

Show $\mathcal{L}(0) = \left(\frac{\partial}{\partial \epsilon}\right)^2: X \rightarrow Y$ Fredholm

(a) $\ker \mathcal{L}(0) = \{\text{constants}\} = \langle 1 \rangle$

indeed $\mathcal{L}(0) \mathcal{L}(0) = 0 \Rightarrow \mathcal{L}(0) = a\epsilon + b$ but $\mathcal{L}(0) \in X \Rightarrow a = \mathcal{L}(0)'(0) = 0$

and $\mathcal{L}(0) = b$ is in $\ker$. Take $\ker \mathcal{L}(0)^\perp := \langle 1 \rangle^\perp$

$$= \left\{ \eta \mid \int_0^1 1 \cdot \eta(\epsilon) d\epsilon = 0 \right\}$$

L^2 orthogonal complement

(b) $g \in \text{Range } \mathcal{L}(0) \Leftrightarrow \exists \mathcal{L} \in X$ s.t. $\mathcal{L}''(\epsilon) = g(\epsilon)$, $g \in C^0$

$$\Rightarrow \mathcal{L}'(\epsilon) = \underbrace{\mathcal{L}'(0)}_{=0} + \int_0^\epsilon g(\epsilon) d\epsilon$$

$$\Rightarrow \mathcal{L}(\epsilon) = \underbrace{\mathcal{L}(0)}_{\text{constant } \ker \mathcal{L}(0)} + \int_0^{\epsilon_1} \left(\int_0^{\epsilon_2} g(\epsilon_2) d\epsilon_2 \right) d\epsilon_1$$

However $\mathcal{L} \in X$ only if $0 = \mathcal{L}'(1) = \int_0^1 g(\epsilon_1) d\epsilon_1 = \langle 1, g \rangle_{L^2}$

Together, rereading the above proves

$$\text{Range } \mathcal{L}(0) = \langle 1 \rangle^\perp \text{ in } C^0 = Y$$

Therefore $\mathcal{L}(0): X \rightarrow Y$ is Fredholm of index = 0

Therefore $\mathcal{L}(\lambda): X \rightarrow Y$ is Fredholm of index = 0

Now check assumption (i), (ii) of C & R

(i) To obtain nontrivial kernel of $\mathcal{L}(\lambda)$ we need

nontrivial solution $\mathcal{L}$ of $\mathcal{L}''(\epsilon) + \lambda \mathcal{L}(\epsilon) = 0$, $\mathcal{L}'(0) = \mathcal{L}'(1) = 0$
for given $\lambda > 0$. ode & $\mathcal{L}'(0) = 0 \Rightarrow \mathcal{L}(\epsilon) = A \cos(\sqrt{\lambda} \epsilon)$

Therefore $\dim \mathcal{L}(\lambda) \leq 1$.

But to have nontrivial kernel element $\mathcal{L} \in X$ we also

require: $0 = \mathcal{L}'(1) = -\sqrt{\lambda} \sin(\sqrt{\lambda})$, i.e. $\lambda = \lambda_k = (k\pi)^2$, $k=1,2,3, \dots$

These $(\lambda_k, 0)$ are candidates for bifurcation points

Let us check (ii),

$\forall g \in Y: Qg = \text{projection onto } (\text{Range } \mathcal{L}(\lambda_k))^\perp$

$$\ker \mathcal{L}(\lambda_k) = \langle \cos(\sqrt{\lambda_k} \epsilon) \rangle = (\text{Range } \mathcal{L}(\lambda_k))^\perp$$

↑
CHECK

(self adjointness of $\mathcal{L}(\lambda)$)

Dyn3 $Q = \text{orthog. proj. onto } \cos(\sqrt{\lambda_k} \xi), \text{ i.e.}$

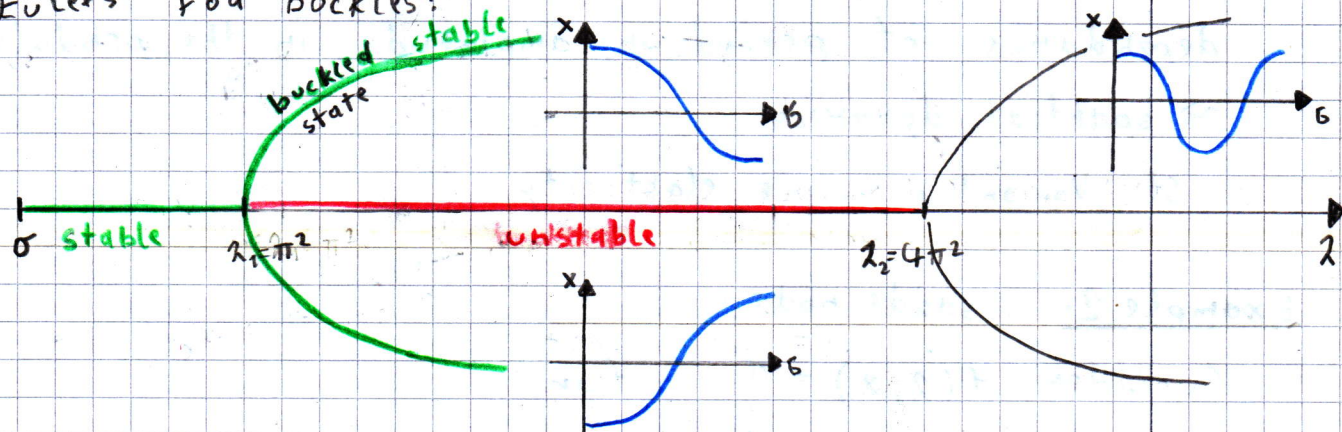
$$Qg = 2 \left(\int_0^1 (\cos(k\pi \xi)) g(\xi) d\xi \right) \cdot \cos(k\pi \xi)$$

$$\Rightarrow Q \mathbb{D}_\lambda \mathbb{D}_x f(0, \xi) \underbrace{\cos(k\pi \xi)}_{\xi \in \text{Ker } \mathbb{L}(\lambda)} = Q \mathbb{D}_\lambda (\xi'' + \lambda \xi) = Q \xi''$$

$$= 2 \left(\int_0^1 \cos(k\pi \xi) \cdot \cos(k\pi \xi) d\xi \right) \cos(k\pi \xi) \\ = 1 \cos(k\pi \xi) \neq 0$$

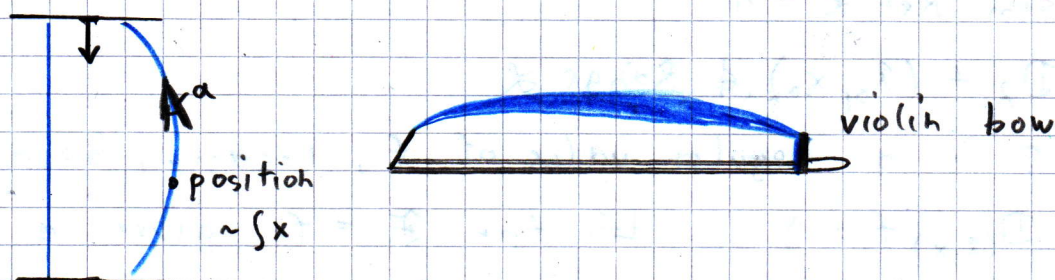
Now C & R show us how:

Euler's rod buckles:



Since $f(\lambda) = -x = -f(\lambda, x)$, the bifurcation branches are not transcritical but symmetric to the λ -axis. Further calculation exhibits supercritical pitchforks.

Buckled rod should look like:

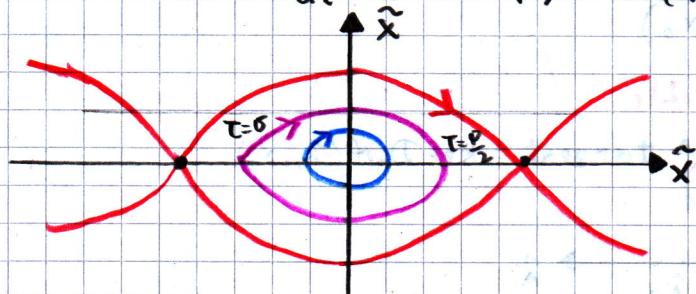


Alternatively, we could view $\odot$ as a pendulum

Let $\tilde{x}(\sqrt{\lambda} \xi) := x(\xi)$. Then

$$\ddot{\tilde{x}} \cdot \lambda + \lambda \sin \tilde{x} = x'' + \lambda \sin x = 0$$

$$\text{with } \cdot = \frac{d}{dt} \quad \tilde{x}(0) = \tilde{x}(\sqrt{\lambda}) = 0 \quad \ddot{\tilde{x}} + \sin \tilde{x} = 0$$



Therefore solutions of $\tilde{x}$ with minimal period p

provide solutions x

$$\sqrt{\lambda} = k \cdot \frac{p}{2} \quad \text{i.e.: } \lambda = k^2 \frac{1}{4} p^2$$

p : as a function of amplitude is called the period map ("time map") [$\leadsto$ Carlos Rocha].

For amplitude $\downarrow 0$ we have $p \rightarrow 2\pi$, hence $\lambda \rightarrow (k\pi)^2 \lambda_k$

The dependence $\lambda(\sigma)$, $x(\sigma)$ of the bifurcating Euler branches are all related to one and the same dependence of period on amplitude, in the pendulum $\leadsto$ spatial dynamics

St. Variant principle elasticity

Example 2: saddle node

Consider $f(\lambda, x) = 0$ for
 $f: \mathbb{R} \times X \rightarrow Y \quad C^k, \quad k \geq 1$

Do not assume any "trivial branch". Instead suppose

$$f(\lambda_0, x_0) = 0$$

$$D_x f(\lambda_0, x_0) = \mathcal{L} = \text{Fredholm, index} = 0$$

i) $\dim \text{Ker } \mathcal{L} = 1$

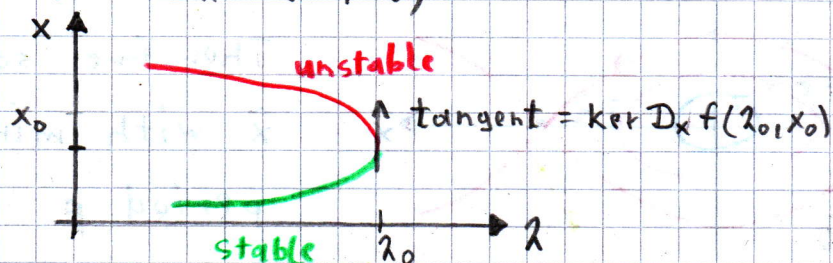
ii) $D_\lambda f(\lambda_0, x_0) \notin \text{Range } \mathcal{L}$

Then 0 is a regular value of f , locally, because
 $\text{Range } D_{(\lambda, x)} f = Y$ L.S. for $\mathcal{F} := f$ gives a
 0 -dim L.S. reduced equation $\phi = 0$ i.e., nothing
to solve there.

Solutions are therefore given by $v = V(u)$ alone,

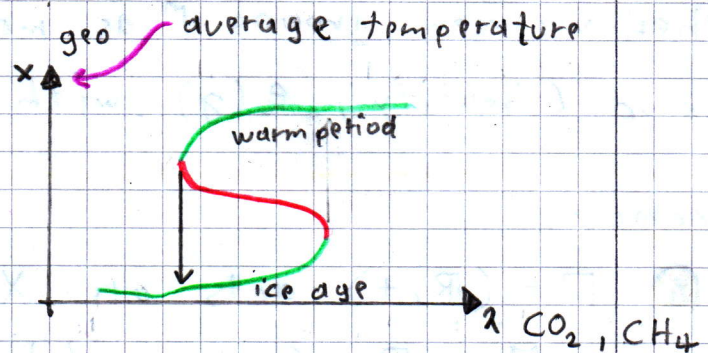
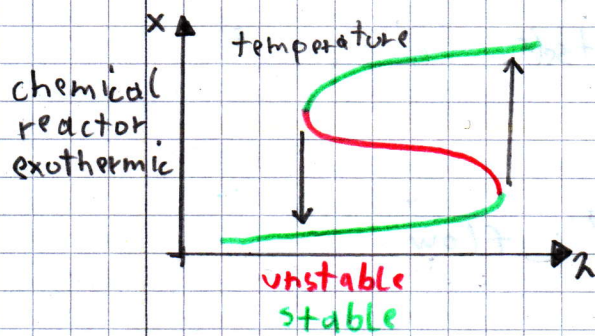
i.e. $\lambda = \lambda(s)$, $x = x(s)$ and $\lambda'(0) = 0$, $\lambda_0 = \lambda(0)$, $x_0 = x(0)$,

$$x = x(s) \in \text{Ker } D_x f(\lambda_0, x_0)$$



For $\dot{x} = f(\lambda, x) \in X = \mathbb{R}^N$, $Y = \mathbb{R}^N$ we can invoke cmf again and obtain, e.g., stability analysis.

Global examples:



hysteresis. slow increase / decrease of λ on lower / upper branch.

7.3 Stationary symmetry breaking

Let Γ be a group which acts on a set X ; i.e. we have a map

$$\begin{aligned} \rho: \Gamma \times X &\rightarrow X & \rho(\text{id}) &= \text{id} \\ (\gamma, x) &\mapsto \rho(\gamma, x) = \rho(\gamma)x = \gamma x & \rho(\gamma^{-1}) &= \rho(\gamma)^{-1} \end{aligned}$$

such that $\rho(\gamma_1) \circ \rho(\gamma_2) = \rho(\gamma_1 \circ \gamma_2)$

For any $x \in X$ we can define the isotropy (stabiliser, fix group, "Standgruppe")

$$\Gamma_x := \{\gamma x = x\} \subseteq \Gamma$$

subgroup

Conversely, given any subgroup K of Γ , we have the K -fixed set (place, space)

$$X_K := \{x \in X \mid \gamma x = x, \text{ for all } \gamma \in K\} = \text{Fix}(K)$$

Note: $x \in X_K \Leftrightarrow \Gamma_x \supseteq K$

We call ρ a representation on X , Banach if all $\rho(\gamma)$ are bounded linear: $X \rightarrow X$

For $X = \mathbb{R}^N$ we have (invertible) $N \times N$ matrices

$\rho(\gamma)$ and hence a group homomorphism

$$\rho: \Gamma \rightarrow GL(N)$$

$$\gamma \mapsto \rho(\gamma)$$

When we view groups Γ as matrix group, we automatically mix up (identify) $\rho(\gamma)$ with γ itself.

Example:

① $\Gamma = (\mathbb{R}, +)$ acts on $X = \mathbb{R}^N$: flow

$$\Gamma_x = \mathbb{R} \text{ for equilibria}$$

$$= p\mathbb{Z} \text{ for periodic orbits}$$

② Γ acts on $X = \Gamma$ by conjugation $\rho(\gamma)g = \gamma g \gamma^{-1}$

$$\text{Then } \Gamma_g = \{ \gamma \mid \gamma g = g \gamma \} = \text{centralizer}$$

$$X_k = \{ g \in \Gamma \mid \gamma g = g \gamma \} = ?$$

③ Γ acts on $X = \{ \text{subgroups of } \Gamma \}$ by

$$\text{conjugation. Then } \Gamma_K = \{ \gamma \mid \gamma K \gamma^{-1} = K \} = \text{normalizer of } K$$

Let $\mathcal{F}: W \rightarrow Z$, C^k , Banach, $\mathcal{F}(\sigma) = \sigma$, $k \geq 1$.

Let ρ_W, ρ_Z be representations of the group Γ

on W, Z . We call $\mathcal{F}$ equivariant w.r. to Γ if

$$\mathcal{F}(\rho_W(\gamma)w) = \rho_Z(\gamma)\mathcal{F}(w) \quad \forall w \in W, \gamma \in \Gamma$$

$$\text{short: } \mathcal{F}(\gamma, w) = \gamma \mathcal{F}(w)$$

[Vanderbauwhede] [Golubitsky & Stewart]

[Chossat & Lauterbach]

Theorem (equivariant LS reduction)

Assume $\mathcal{L} := \mathcal{F}(\sigma)$ is quasi-Fredholm with projections

$P; Q = \text{id} - P$ in Z onto $(\text{Range } \mathcal{L})^{(c)}$ and projections

$\tilde{P}, \tilde{Q} = \text{id} - \tilde{P}$ in W onto $(\text{Ker } \mathcal{L})^{(c)}$

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(before: $w = u \oplus v = (u, v)$; now: $u = \tilde{P}w, v = \tilde{Q}w$)

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Then $F(u, v) = \sigma$, (locally) $\Leftrightarrow \begin{cases} PF(u, v) = \sigma \Leftrightarrow v = V(u) \\ \text{and } \Phi(u) := QF(u, V(u)) = \sigma \end{cases}$ L-S-reduced equation; $\Phi, V \in C^k, V(\sigma) = \sigma, \Phi(\sigma) = \sigma, V'(\sigma) = \sigma$ $\Phi'(\sigma) = \sigma$

AND

 $\Phi: \tilde{P}W \rightarrow QZ$ are both
equivariantRemark: $P: Z \rightarrow Z$ equivariant $\Rightarrow \text{Ker } P, \text{Range } P$ are Γ -invariant. Indeed: $\gamma Pz \stackrel{\text{equiv.}}{=} P\gamma z$ Similarly $Pz = \sigma$ $\text{Range } P \quad \text{Range } P$ $\Rightarrow \sigma = \gamma Pz = P\gamma z \Rightarrow \gamma z \in \text{Ker } P$. Therefore the restrictions of $R_W(\Gamma), R_Z(\Gamma)$ provide representations on $u, v, \text{Range } \mathcal{L}$.

(Proof on page 44)

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continuous reps $\mathcal{S}: \Gamma$ topological group

(i.e. composition is conto)

 $X = \text{Banach}, \mathcal{S}(\gamma)$ bd. linear, $\gamma \mapsto \mathcal{S}(\gamma) \in \mathcal{L}(X)$ cont.

(reps are assumed to be cont., wherever necessary)

Note: $X_K = \bigcap_{\gamma \in K} \text{Ker}(\mathcal{S}(\gamma) - \text{id}) = \text{closed subspace} \Rightarrow \text{also Banach}$ Equivariance: $F: W \rightarrow Z$ equivariant ("invariant", "covariant")if $F(\mathcal{S}_W(\gamma)w) = \mathcal{S}_Z(\gamma)F(w) \quad (\forall w, \gamma)$ Rep.: Let $\mathcal{L}: W \rightarrow Z$ be linear and equivariant.Then $\text{Ker } \mathcal{L}, \text{Range } \mathcal{L}$ are Γ -invariant i.e. $u \in \text{Ker } \mathcal{L}$ $\Rightarrow \mathcal{S}(\gamma)u \in \text{Ker } \mathcal{L} \quad (\forall \gamma)$ Proof: $\mathcal{L} \mathcal{S}(\gamma)u = \mathcal{S}(\gamma)\mathcal{L}u = \sigma$ for $u \in \text{Ker } \mathcal{L}$ Similarly for Range $\square$ Corollary: Eigenspaces and generalized eigenspaces of equivariant operators $\mathcal{L}$ are invariant.Proof: Consider $\mathcal{L} - \lambda \text{id}$, instead of $\mathcal{L}$.Example: Consider $\dot{x} = f(x), X = \mathbb{R}^N$, ode. Thenequivariance $f(\gamma x) = \gamma f(x)$ for all $\gamma \in \Gamma$, group $x \in X$

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